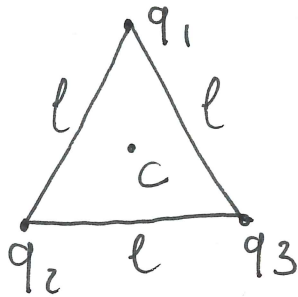


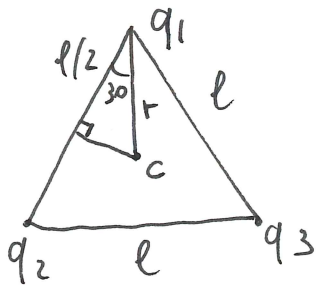
# Exercise Session 1

-1-

Calculate the electrostatic potential in C.



Solution.



$V_1 = \frac{q_1}{4\pi\epsilon_0 r}$  potential in C due to charge  $q_1$ ,

where  $r = \frac{l}{\sqrt{3}}$ , so  $V_1 = \frac{\sqrt{3}}{4\pi\epsilon_0 l} \cdot q_1$

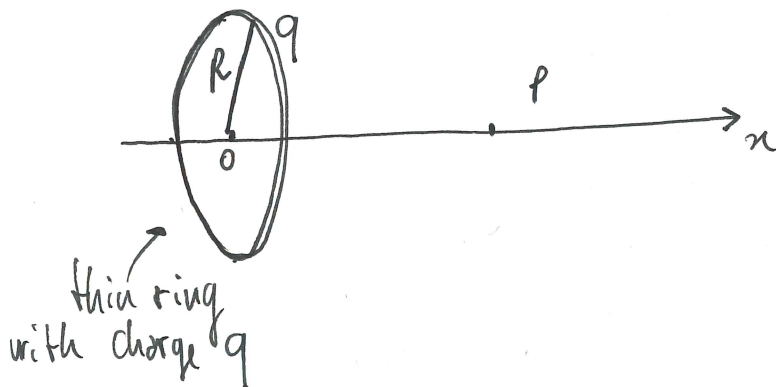
Thanks to the superposition rule, the total potential in C is given by the algebraic sum of all contributions from each single charge.

$$V = \sum_{i=1}^3 V_i = \frac{\sqrt{3}}{4\pi\epsilon_0 l} \sum_{i=1}^3 q_i$$

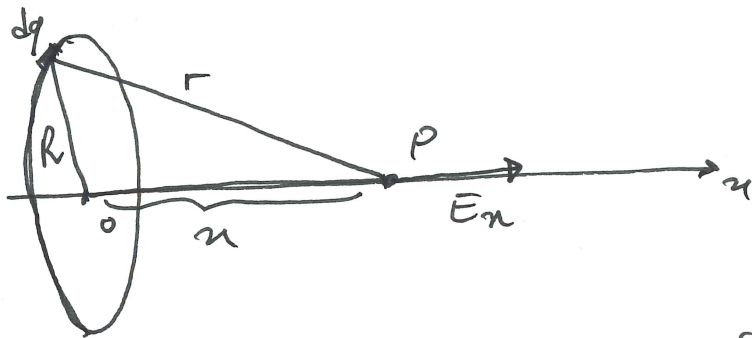
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-2-

Calculate the electrostatic field  $\vec{E}$  and potential  $V$  at an arbitrary point P along the x-axis. The thin ring has radius  $R$  and a total charge  $q$ .



Solution



The charge distribution is one-dimensional so that a linear charge density can be defined as  $\lambda = \frac{q}{2\pi R} \frac{C}{m}$

Each infinitesimal element  $dl$  holds a charge  $dq = \lambda \cdot dl$ .  
 Each element  $dl$  is treated as a point-like source, producing an infinitesimal potential  $dV$  at a distance  $r$ .

$$dV = \frac{dq}{4\pi\epsilon_0 r} = \frac{\lambda dl}{4\pi\epsilon_0 r} \quad \text{where } r = \sqrt{R^2 + x^2}$$

$$dl = R dp$$



After integration:

$$V = \int_0^{2\pi} \frac{\lambda R dp}{4\pi\epsilon_0 r} = \frac{q}{4\pi\epsilon_0 \sqrt{R^2 + x^2}}$$

If  $x \gg R$ ,  $\sqrt{R^2 + x^2} = x \cdot \sqrt{1 + \frac{R^2}{x^2}} \sim x \left(1 + \frac{1}{2} \frac{R^2}{x^2} + \dots\right) \sim x$   
 that is: the potential is the same as the charge  $q$  would be concentrated in  $O$ .

\* note: The integration constant  $C=0$  because this expression for  $dV$  is equivalent to the potential of a single point-like charge.

The electrostatic field is calculated upon derivation of  $V$ .

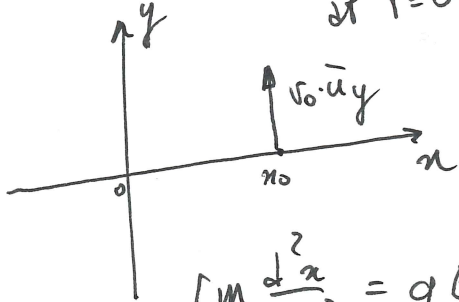
$$\begin{cases} E_x = -\frac{\partial V}{\partial x} = -\frac{q}{4\pi\epsilon_0} \frac{\partial}{\partial x} \left( \frac{1}{\sqrt{x^2 + R^2}} \right) = \frac{q}{4\pi\epsilon_0} \frac{x}{(R^2 + x^2)^{3/2}} \\ E_y = -\frac{\partial V}{\partial y} = 0 \\ E_z = -\frac{\partial V}{\partial z} = 0 \end{cases}$$

$\vec{E}$  is oriented parallel to the  $x$ -axis.

-3-

Express the equation of motion for a particle with electric charge  $q$  and mass  $m$ , such that  $\frac{q}{m} = 5 \cdot 10^6 \frac{C}{kg}$ , within an electrostatic potential defined over the  $xy$  plane as  $V(x, y) = V_0(x^2 + y^2)$ , where  $V_0 = 10^7 \frac{V}{m^2}$ . At time  $t=0$ , the particle is in point  $P(x_0, 0)$  and has a velocity  $\vec{v} = v_0 \vec{u}_y$ , where  $x_0 = 1 \text{ cm}$  and  $v_0 = 10^5 \frac{m}{s}$ .

Solution:



$$\vec{E} = -\nabla V = \vec{u}_x \left( -\frac{\partial V}{\partial x} \right) + \vec{u}_y \left( -\frac{\partial V}{\partial y} \right)$$

$$\vec{E} = -2V_0 x \cdot \vec{u}_x - 2V_0 y \cdot \vec{u}_y$$

From the Newton's 2nd law:

$$\begin{cases} m \frac{d^2 x}{dt^2} = q E_x = -2V_0 x \cdot q & \text{motion along } x \\ m \frac{d^2 y}{dt^2} = q E_y = -2V_0 y \cdot q & \text{motion along } y \end{cases}$$

The two equations (differential) allow for harmonic solutions:

$$x(t) = x_0 \cos \omega t \quad \text{where } x_0 \text{ is known } (x(0) = x_0)$$

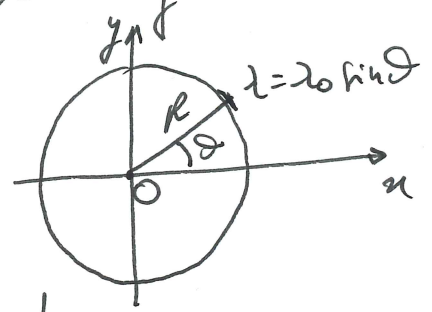
$$y(t) = y_0 \sin \omega t \quad \text{where } y_0 \text{ must be calculated using the information on the velocity.}$$

$$v_y(t) = \frac{dy}{dt} = \omega y_0 \cos \omega t. \quad \text{At } t=0, v_y(0) = v_0, \text{ therefore: } \omega y_0 = v_0$$

$$y(t) = \frac{v_0}{\omega} \sin \omega t.$$

-4-

Calculate the expression of the electrostatic field in the center O of an annular distribution of electric charge having linear density  $\lambda = \lambda_0 \sin \theta$ , as depicted in the figure -



Solution.

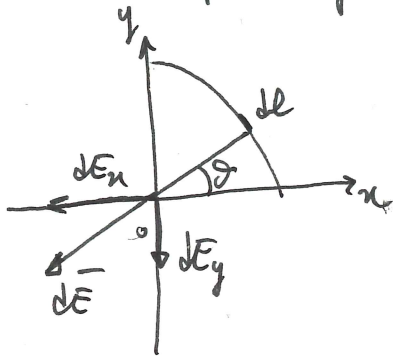
The total electric field in the center is given by the sum (integral) of all the infinitesimal contributions  $d\vec{E}$  from corresponding infinitesimal arcs  $dl$  on the ring.

Each arc  $dl$  brings an infinitesimal charge

$$dq = \lambda \cdot dl = \lambda_0 \sin \theta \cdot dl$$

$$dE_x = - \frac{dq}{4\pi\epsilon_0 R^2} \cos \theta = - \frac{\lambda_0 \sin \theta \cdot dl}{4\pi\epsilon_0 R^2} \cos \theta$$

$$dE_y = - \frac{dq}{4\pi\epsilon_0 R^2} \sin \theta = - \frac{\lambda_0 \sin \theta \cdot dl}{4\pi\epsilon_0 R^2} \sin \theta$$



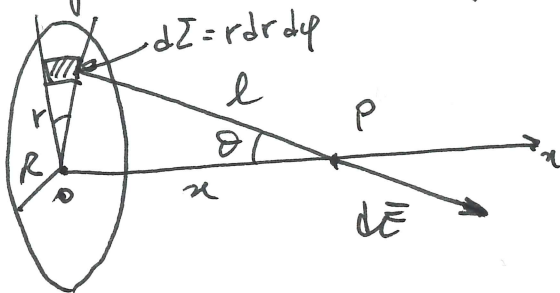
Taking in mind that  $dl = R d\theta$

$$E_x = - \frac{\lambda_0}{4\pi\epsilon_0 R} \int_0^{2\pi} \sin \theta \cdot \cos \theta \, d\theta = 0$$

$$E_y = - \frac{\lambda_0}{4\pi\epsilon_0 R} \int_0^{2\pi} \sin^2 \theta \, d\theta = - \frac{\lambda_0}{4\epsilon_0 R}$$

-5-

Express the electrostatic field produced by a disk-shaped planar distribution of electric charge with constant density  $\sigma$  ( $\frac{C}{m^2}$ ) at an arbitrary distance  $x$  along its axis.



Solution

Each infinitesimal surface  $dZ$  contributes with a field  $d\vec{E}$  to the total field in P.

For symmetry reasons, the field  $\vec{E}$  will only have  $x$ -component as all the transverse components will be cancelled out pairwise —

$$dE_x = \frac{dq}{4\pi\epsilon_0 l^2} \cos\theta \quad \text{where} \quad dq = \sigma \cdot dZ = \sigma \cdot r dr dp$$
$$l = \sqrt{x^2 + r^2}$$
$$\cos\theta = \frac{x}{\sqrt{x^2 + r^2}}$$

Integrating on  $r$  and  $p$  leads to:

$$E_x = \int_0^{2\pi} dp \int_0^R dr \frac{\sigma x}{4\pi\epsilon_0} \frac{r}{(x^2 + r^2)^{3/2}} = \frac{\sigma x}{2\epsilon_0} \left( \frac{1}{x} - \frac{1}{\sqrt{x^2 + R^2}} \right)$$

Alternative approach: Calculate  $E_x = -\frac{\partial V}{\partial x}$ .

$$dV = \frac{dq}{4\pi\epsilon_0 l} = \frac{\sigma r dr dp}{4\pi\epsilon_0 \sqrt{x^2 + r^2}}$$

$$V = \frac{\sigma}{4\pi\epsilon_0} \int_0^{2\pi} dp \int_0^R dr \frac{r}{\sqrt{x^2 + r^2}} = \frac{\sigma}{2\epsilon_0} \left[ \sqrt{x^2 + r^2} \right]_0^R = \frac{\sigma}{2\epsilon_0} (\sqrt{x^2 + R^2} - x)$$

$$E_x = -\frac{\partial V}{\partial x} = \frac{\sigma}{2\epsilon_0} - \frac{\sigma x}{2\epsilon_0 \sqrt{x^2 + R^2}}$$