

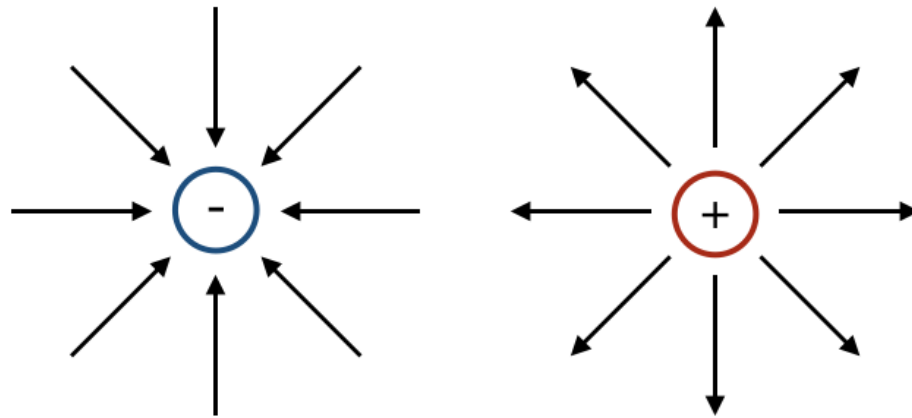
Lecture 3: stationary electric field

- 1) Energy in electric field
- 2) Boundary condition in electro-statics
- 3) perfect conductor
- 4) Ohm's law

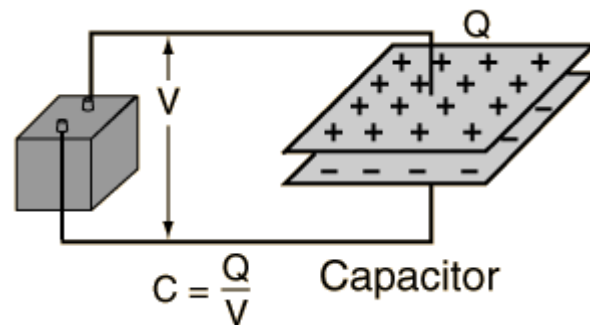
Electric field and energy

Electric potential energy : energy required to move a charge through an electric field.

$$E_p = \frac{qQ}{4\pi\epsilon_0 r}$$



Capacitors are devices that can store electric potential energy and release it as charge through an electric circuit.



Electric energy in capacitor

The energy stored on a capacitor can be expressed in terms of the work done by the battery.

Voltage represents energy per unit charge

$$W_e = \int_a^b F dl = \int_a^b qE dl = qV$$

The work to move a charge element dq from the negative plate to the positive plate is equal to $V dq$

$$dW_e = V dq$$

↓

$$W_e = \int_0^Q V dq = \int_0^Q \frac{q}{C} dq = \frac{Q^2}{2C} = \frac{1}{2} CV^2$$

↓

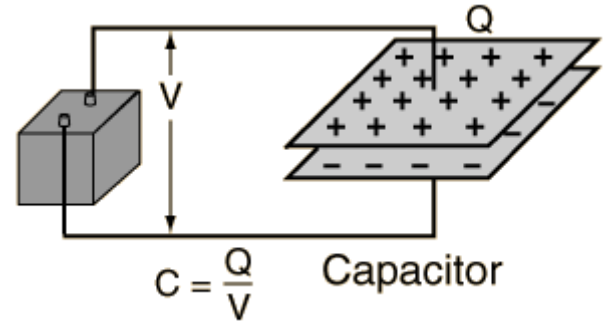
$Q = CV$

$$W_e = \int IU dt = \int_0^V C \frac{dU}{dt} U dt = \int_0^V CU dU = \frac{1}{2} CV^2$$

←

$$I = C \frac{dU}{dt}$$

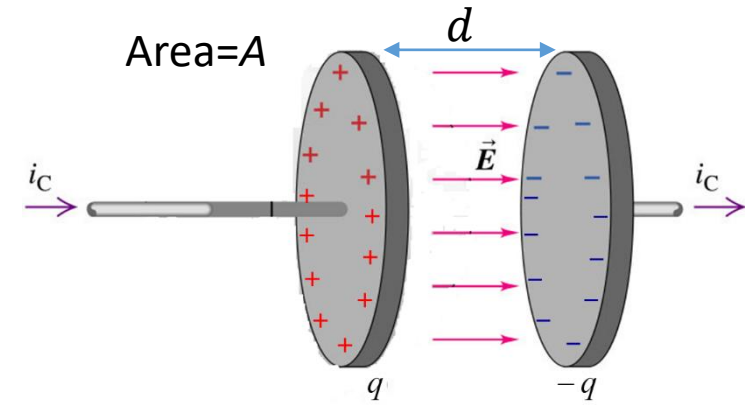
$$W_e = \frac{1}{2} CV^2$$



Energy in electric field

$$W_e = \frac{1}{2} CV^2 = \frac{1}{2} \frac{\epsilon A}{d} V^2 = \frac{1}{2} \epsilon Ad \frac{V^2}{d^2} = \frac{1}{2} \epsilon V_{vol} E^2 = \frac{1}{2} V_{vol} \mathbf{DE}$$

$C = \frac{\epsilon A}{d}$
 Volume = Ad
 $\frac{V^2}{d^2} = E^2$
 $\mathbf{D} = \epsilon \mathbf{E}$



Electric Energy density $\eta_e = \frac{W_e}{V_{vol}} = \frac{1}{2} \epsilon E^2 = \frac{1}{2} \mathbf{DE}$

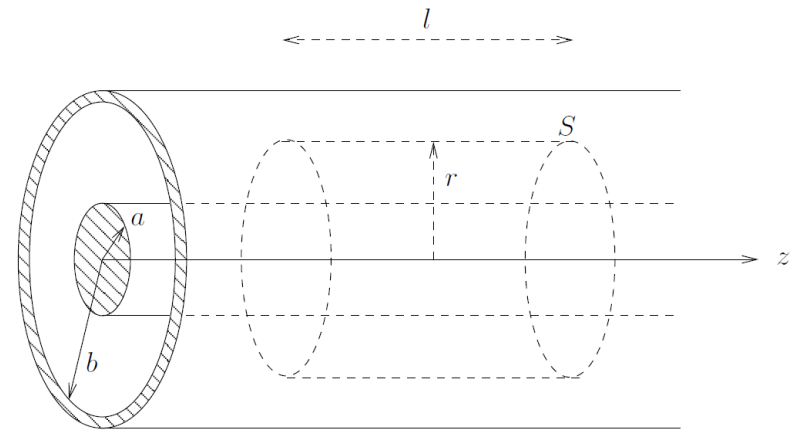
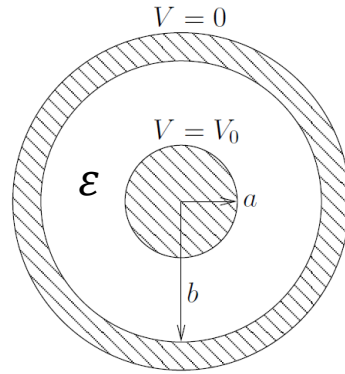
Electric Energy density

$$\eta_e = \frac{1}{2} ED = \frac{1}{2} \epsilon E^2$$

Example:

Assuming : infinite long cable and the permittivity of the dielectric in between is ϵ

- 1) calculate the capacitance per unit length.
- 2) calculate the electric energy stored in unit length



$$C' = \frac{Q'}{V_0} \quad V_0 = \frac{Q'}{2\pi\epsilon} \ln \frac{b}{a} \quad C' = \frac{Q'}{V_0} = \frac{2\pi\epsilon}{\ln \frac{b}{a}}$$

$$W'_e = \frac{1}{2} C' V_0^2 = \frac{\pi\epsilon}{\ln \frac{b}{a}} V_0^2$$

$$W_e = \frac{1}{2} \int_{\text{mellom lederne}} \epsilon E^2 dv = \frac{\epsilon V_0^2}{2 \left(\ln \frac{b}{a}\right)^2} \int_a^b \frac{2\pi r dr l}{r^2} = \frac{2\pi\epsilon V_0^2}{2 \ln \frac{b}{a}} l$$

$$W'_e = \frac{W_e}{l} = \frac{\pi\epsilon}{\ln \frac{b}{a}} V_0^2$$

Boundary conditions in electrostatics

Electric field involves more than one materials

Normal component

$$D_{n1}\Delta S - D_{n2}\Delta S = \rho_s \Delta S \quad \oint \mathbf{D} \cdot d\mathbf{S} = Q_s$$

$$D_{n1} - D_{n2} = \rho_s$$

Conservative field: electrostatic electric field

Tangential component

$$E_{t1}\Delta l - E_{t2}\Delta l = 0 \quad \oint \mathbf{E} \cdot d\mathbf{l} = 0$$

$$E_{t1} = E_{t2}$$

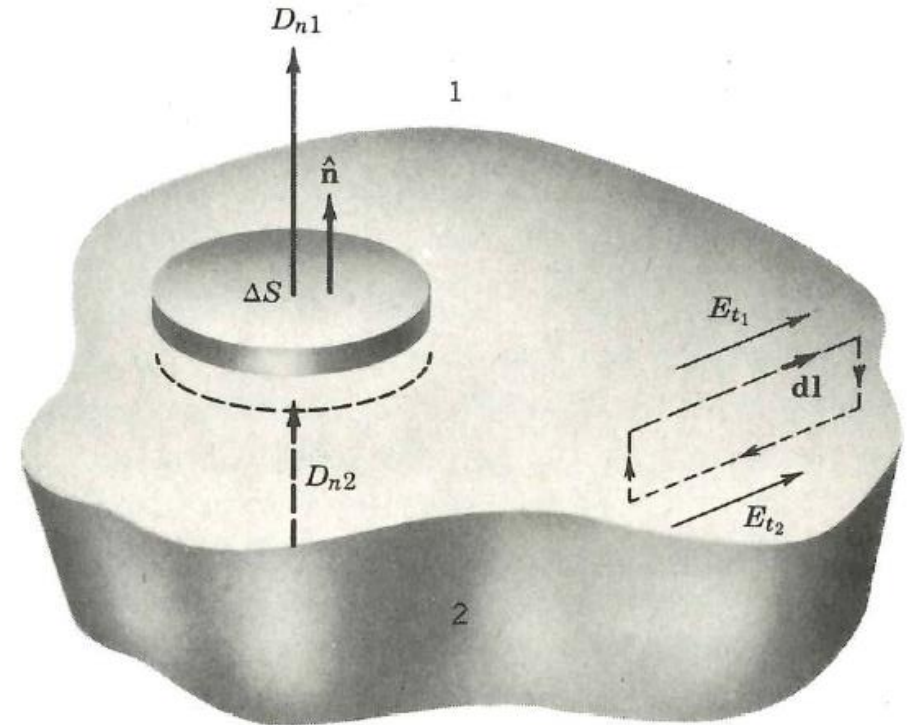


Fig. 1.14a Boundary between two different media.

Electric field changes direction across two different dielectric materials

Considering, no charge on the surface between two dielectric materials

$$D_{n1} - D_{n2} = \rho_s = 0;$$



$$\epsilon_1 E_{n1} = \epsilon_2 E_{n2}$$



$$E_{n2} = \frac{\epsilon_1 E_{n1}}{\epsilon_2}$$

$$E_{t1} = E_{t2}$$



$$E_{t2} = E_{t1} = E_{n1} \tan \theta_1$$

$$\tan \theta_1 = \frac{E_{t1}}{E_{n1}}$$

$$\theta_2 = \tan^{-1} \frac{E_{t2}}{E_{n2}} \Rightarrow \theta_2 = \tan^{-1} \left(\frac{\epsilon_2}{\epsilon_1} \tan \theta_1 \right)$$

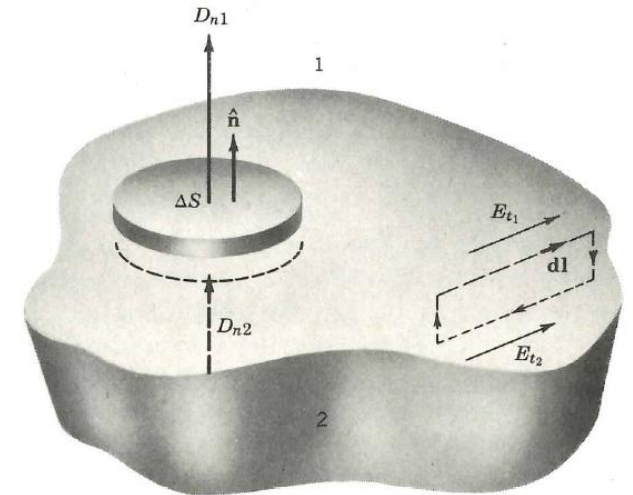
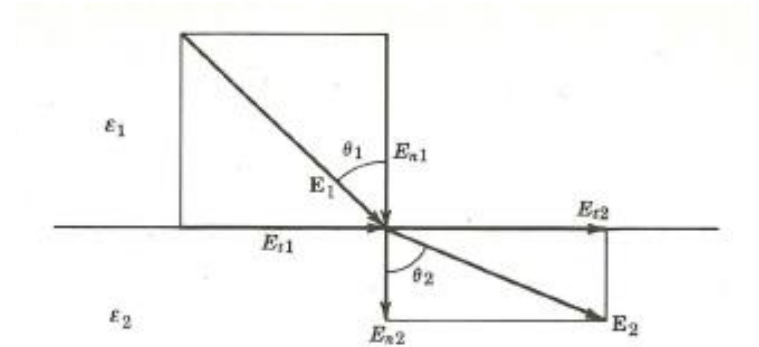


Fig. 1.14a Boundary between two different media.

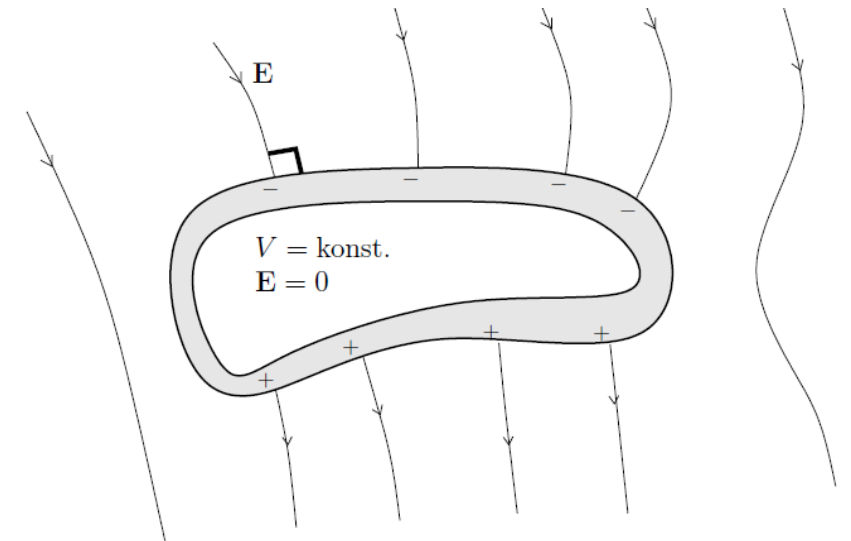
Perfect conductor in electric field

- 1) $E=0$ inside the conductor. Both $E_t = E_n=0$.
- 2) $\rho_{in} = 0$, no charge inside the conductor.
- 3) $\rho_s \neq 0$, there is surface charge .
- 4) The electric field outside the boundary of the perfect conduct is

$$E_n = \frac{\rho_s}{\epsilon}, E_t = 0,$$

- 5) The conductor is an equipotential surface.

$$V_{AB} = - \int_A^B E dl = 0.$$



$$D_{n1} \Delta S - D_{n2} \Delta S = \rho_s \Delta S$$

$$\oint \mathbf{E} \cdot d\mathbf{l} = E_{t1} \Delta l - E_{t2} \Delta l = 0$$

Faraday cage:

$$D_{n1} - D_{n2} = \rho_s$$

$$E_{t1} = E_{t2}$$

<https://www.youtube.com/watch?v=t23iXhEiQUc>

Material conductivity and ohm's law

Good conductive material:

Silver: $6.2 \times 10^7 S/m$

Copper $5.8 \times 10^7 S/m$

Gold $4.1 \times 10^7 S/m$

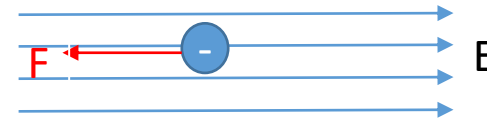
Aluminium $3.5 \times 10^7 S/m$

Non conductive material

Glass $\times 10^{-12} S/m$

Air $\times 10^{-14} S/m$

Once there are free charges in an electric field, the charges can move along the electric field direction.



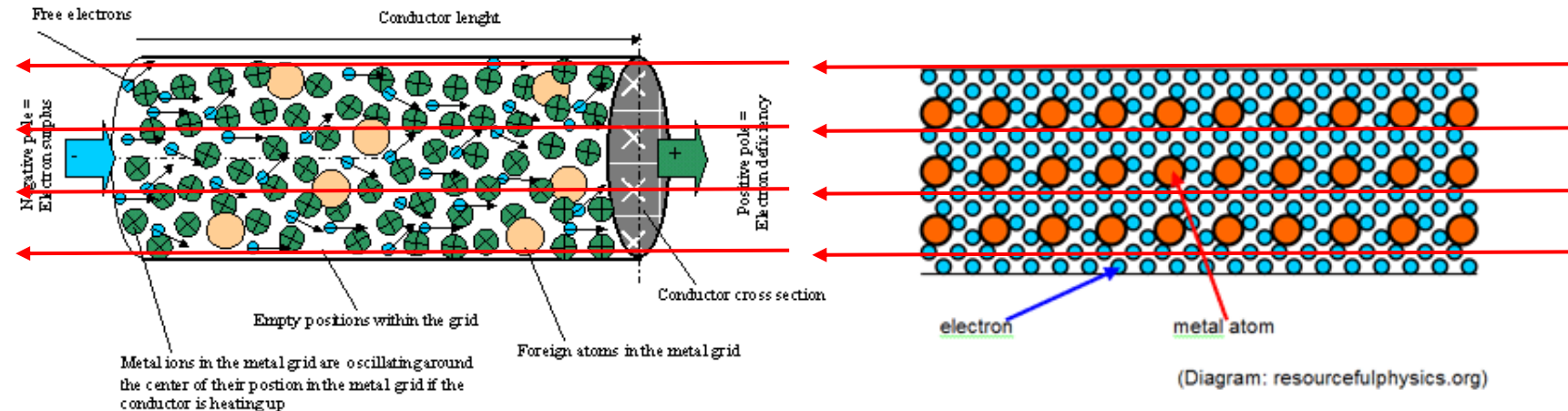
That how many charges can flow through a material is dependent of the material property and the applied electric field.

When electrons move in a material, they collide atoms and lost energy.

The capability is represented by σ , conductivity.

J is called current density. The charges go through the material per unit area per second.

$$J = \sigma E$$



Low conductive material

High conductive material

Ohm's law (II)

An electric current is a flow of electric charge.

$$I = \frac{dQ}{dt}$$

Current density J is the current per unit area

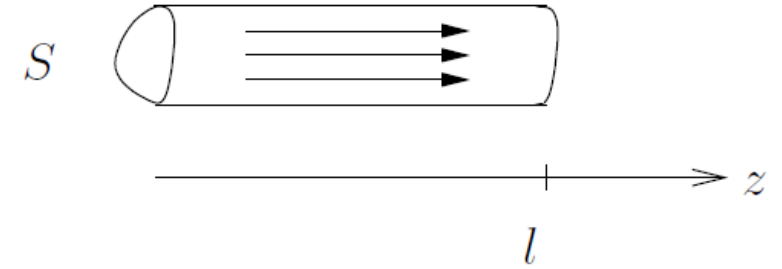
$$I = \int_S \mathbf{J} \cdot d\mathbf{S}.$$

$$\mathbf{J} = \sigma \mathbf{E}$$

$$R = \frac{V}{I}$$

Example:

A conductor with constant conductivity σ , and cross-sectional area S , the length is l and the constant current is I



- 1) Calculating the resistance of the conductor and deriving $R = \frac{V}{I}$.
- 2) Calculating the power done by the current.

Area S , and I is also constant. $J = I/S$

Constant σ , the electric field is constant: $E = J/\sigma$

$$V = \int_0^l E dz = \int_0^l (J/\sigma) dz = Jl/\sigma$$

$$R = \frac{l}{S\sigma} = \frac{Jl}{\sigma JS} = \frac{V}{I}$$

$$\text{Vector: } \vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{r}$$

$$\text{Vector: } \vec{F} = \frac{qQ}{4\pi\epsilon_0 r^2} \hat{r} \quad : \text{Coulomb's law}$$

$$\vec{E} = \vec{F}/q \quad V = - \int \mathbf{E} \cdot d\mathbf{l}$$

$$W_e = \int_0^l F dl = \int_0^l qE dl = qV \quad \Rightarrow \quad W_e = \int VI dt$$

$q = \int I dt$

$$P = VI = I^2 R$$

Example:

A solid conductive ball with a radius a is put into a hollow conductive ball with inner radius b , between is a type material with conductivity σ .

- What is the resistance between the two balls?
- If the solid ball is buried deeply in to earth, what is the earth resistance?

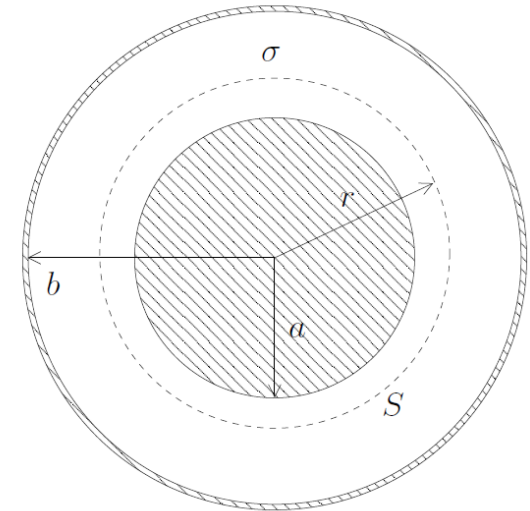
$$a = 0.5 \text{ m} \quad \sigma = 10^{-2} \text{ m}^{-1}\Omega^{-1}$$

$$\mathbf{J} = \frac{I}{4\pi r^2} \hat{\mathbf{r}}, \quad \text{for } a < r < b.$$

$$\mathbf{E} = \mathbf{J}/\sigma = I\hat{\mathbf{r}}/(4\pi\sigma r^2)$$

$$V = -\int_a^b E dr = -\int_a^b \frac{I}{4\pi\sigma r^2} dr = \frac{I}{4\pi\sigma} \left(-\frac{1}{b} + \frac{1}{a} \right)$$

$$R = \frac{V}{I} = \frac{1}{4\pi\sigma} \left(\frac{1}{a} - \frac{1}{b} \right)$$



example

Now if a half solid ball is buried in earth as shown in picture, recalculate the resistance?

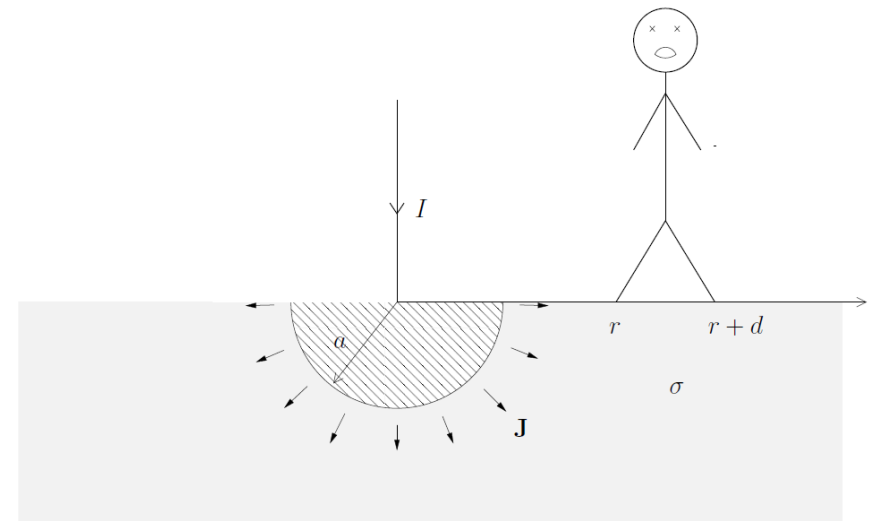
If the current is 1000A $\sigma = 10^{-2} \text{ m}^{-1}\Omega^{-1}$, $r = 1 \text{ m}$ og $d = 0.75 \text{ m}$

What is the voltage between the two legs of the people

$$\mathbf{J} = \frac{I}{2\pi r^2} \hat{\mathbf{r}}, \quad \mathbf{E} = \frac{I}{2\pi\sigma r^2} \hat{\mathbf{r}}, \quad \text{for } r > a.$$

$$R = \frac{1}{2\pi\sigma a}$$

$$V = -\int_r^{r+d} E(r) dr = -\int_r^{r+d} \frac{I}{2\pi\sigma r^2} dr = \frac{I}{2\pi\sigma} \left(\frac{1}{r} - \frac{1}{r+d} \right)$$



Kirchoff's law

Current conservation:

Kirchoff's law: at any point at electric circuit, the sum of the current is zero.

$$\oint_S \mathbf{J} \cdot d\mathbf{S} = 0.$$