

TFE4120 Electromagnetism: crash course

Intensive course: Two-weeks.

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Exercises help:

Participants: should have Bsc in electronic, electrical/ power engineering.

Aim of the course: Give students a minimum pre-requisity to follow a 2-year master program in electronics or electrical /power engineering.

Webpage: All information is posted there .

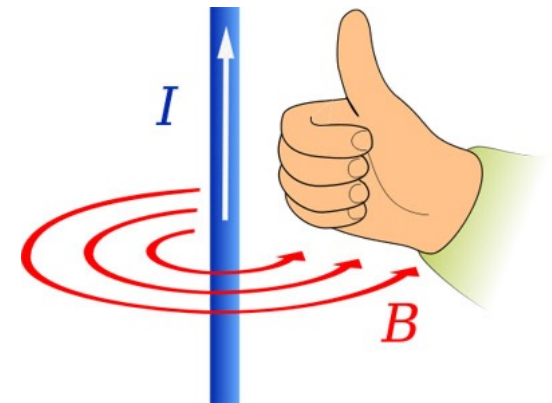
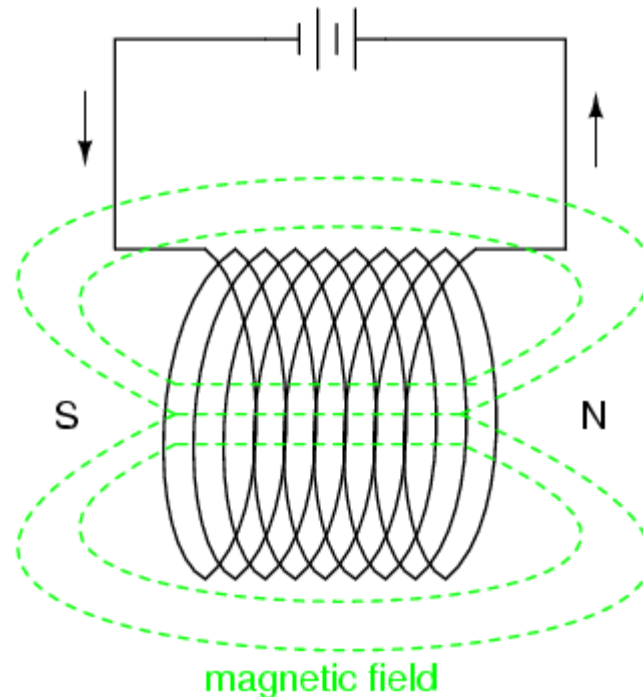
<https://www.ntnu.no/wiki/display/tfe4120/Crash+course+in+Electromagnetics+2022>

Content for lectures:

- Lecture 1: Introduction and Vector calculus
- Lecture 2: Electro-static
- Lecture 3: Electro-dynamic
- Lecture 4: Magnetic-static
- Lecture 5: Electro-magnetic
- Lecture 6: Electro-magnetic wave

Lecture1: Electro-magnetism and vector calculus

- 1) What does electro-magnetism describe?
- 2) Brief induction about Maxwell equations
- 3) Electric force: Coulomb's law
- 4) Vector calculus

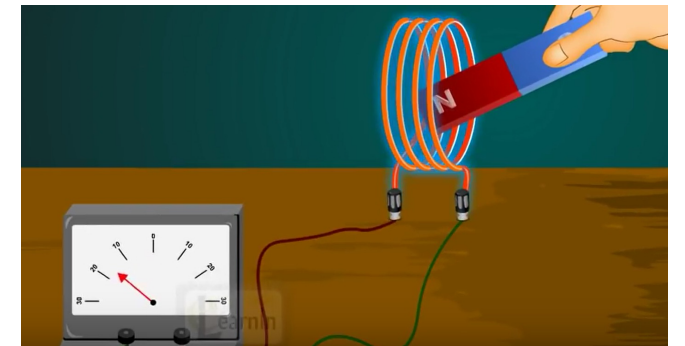
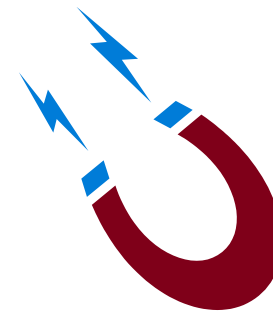
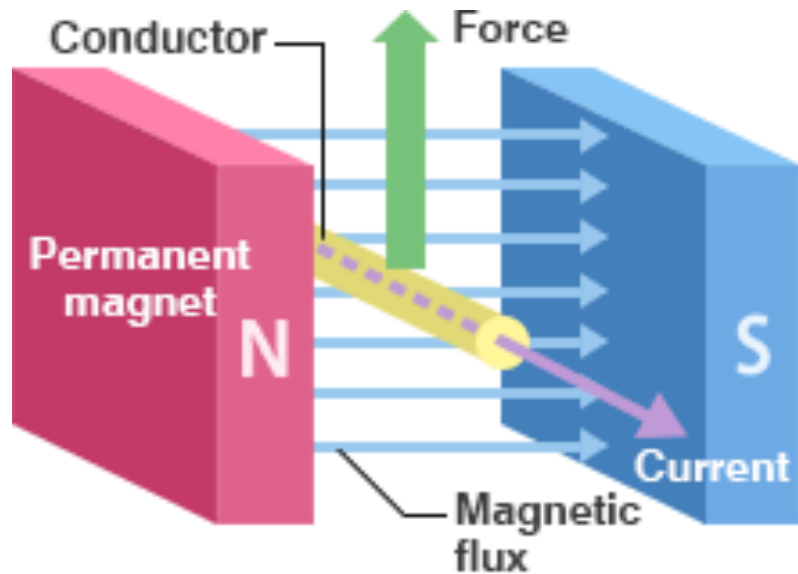


Electro-magnetism

Originally, electricity and magnetism were considered to be two separate phenomena

Electro-magnetism: Physical interaction between electricity and magnetism.

Electro-magnetic force: one of the four fundamental interactions in the nature. (gravitation, electromagnetism, the strong and weak forces)



History:

Carl Friedrich Gauss (1777-1855): German mathematician and physicist

The electric flux out of a closed surface = total enclosed charge divided by the permittivity of free space

Electrostatic

$$\oint \vec{E} \cdot d\vec{A} = \frac{Q_{in}}{\epsilon_0}$$

Andre-Marie Ampere (1775-1836): French physicist and mathematician

The magnetic field produced by an electric current is proportional to the magnitude of the current with a proportionality constant equal to the permeability of free space (μ_0)

Magnetostatic

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I$$

Michael Faraday (1791-1867): English Scientist

- In 1831 Faraday observed that a moving magnet could induce a current in a circuit.
- He also observed that a changing current could, through its magnetic effects, induce a current to flow in another circuit.

Magnetodynamic $V = - \frac{d\phi}{dt}$

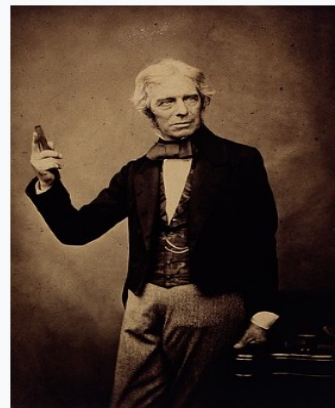
Carl Friedrich Gauss



André-Marie Ampère



Michael Faraday
FRS





Founder of electromagnetism

James Clerk Maxwell: (1839-1879)

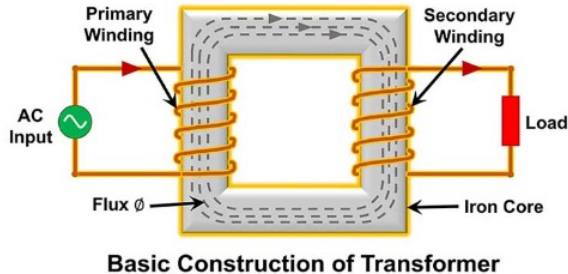
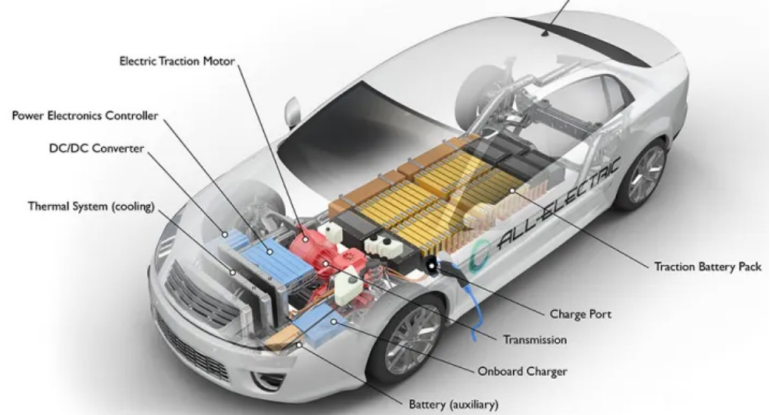
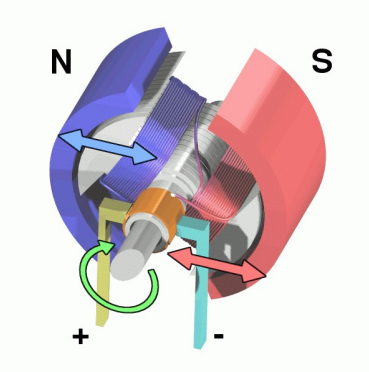
Scottish Mathematician

- Developed a scientific theory to explain electromagnetic waves.
- Coupled the electrical fields and magnetic fields together
- he established the foundations of electricity and magnetism as **electromagnetism**.
- Maxwell equations

Daily life applications

For example:

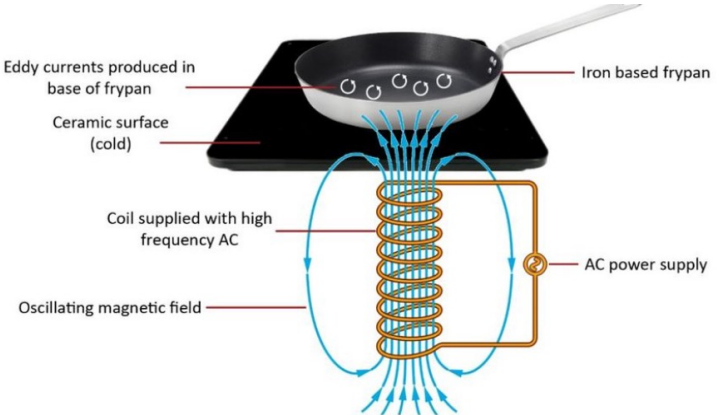
- Electric motor/generator:
- Battery charger
- Induction oven
-



Basic Construction of Transformer

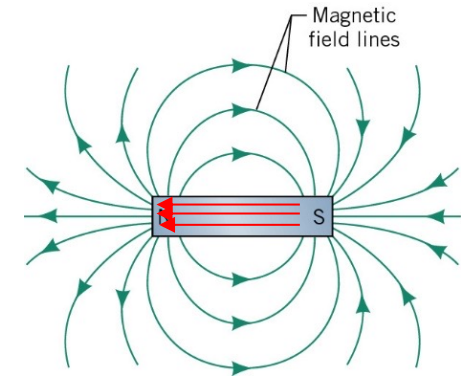
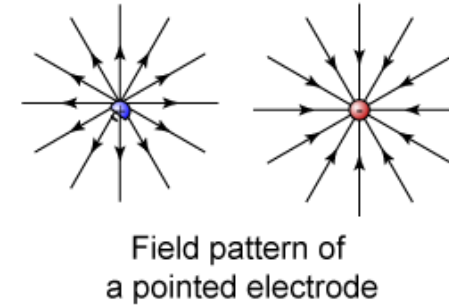


Iphone charger



Electromagnetism: Maxwell equations

- A static **electric charges** produces an electric field
- There is no **magnetic charge** (monopole).
- A **changing magnetic field** produces an electric field,
- **Charges in motion** (an electrical current) produce a magnetic field
- A **changing electric field** produces a magnetic field.



Electric and Magnetic fields can produce forces on charges

$$\nabla \cdot \mathbf{E} = \frac{\rho_v}{\epsilon} \quad \text{(Gauss' Law)}$$

$$\nabla \cdot \mathbf{H} = 0 \quad \text{(Gauss' Law for Magnetism)}$$

$$\nabla \times \mathbf{E} = -\mu \frac{\partial \mathbf{H}}{\partial t} \quad \text{(Faraday's Law)}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \epsilon \frac{\partial \mathbf{E}}{\partial t} \quad \text{(Ampere's Law)}$$

E: electric field, Vector

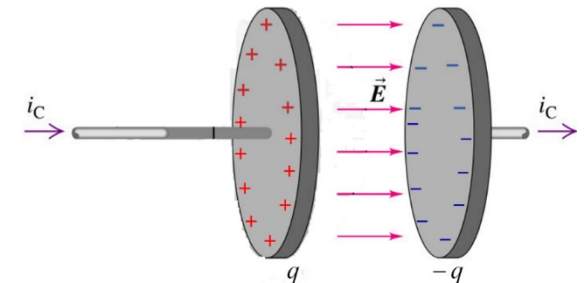
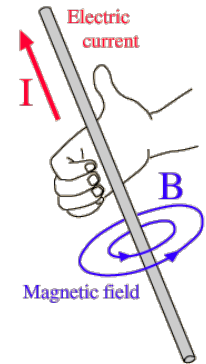
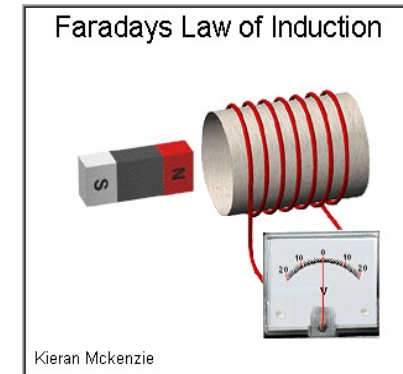
D: electric flux density, Vector

H: magnetic field, Vector

B: magnetic flux density, Vector

J: current density, Vector

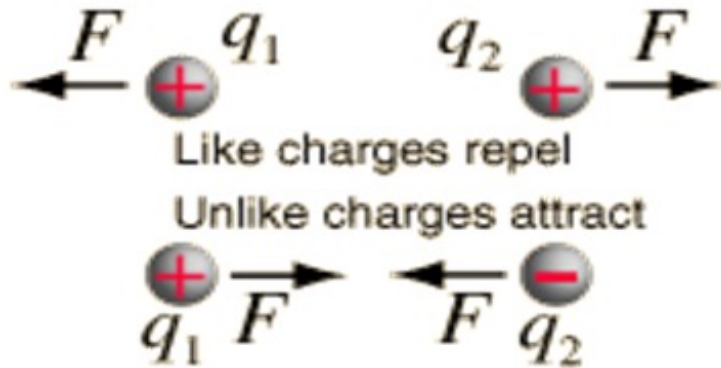
ρ : Static charge density



Electricity and magnetism had been unified into electromagnetism!

Coulomb's law: force between electrostatic charges

Published in 1785 by French physicist [Charles-Augustin de Coulomb](#) and was essential to the development of the [theory of electromagnetism](#)



$$\text{Scalar: } F = k \frac{q_1 q_2}{r_{12}^2} = \frac{q_1 q_2}{4\pi\epsilon_0 r_{12}^2}$$

$$\text{Vector: } \vec{F} = \frac{q_1 q_2}{4\pi\epsilon_0 r_{12}^2} \hat{r}_{12}$$

\hat{r}_{12} is just for direction, its absolute value is 1.

$$k = \frac{1}{4\pi\epsilon_0} \approx 9 \times 10^9 \text{ N} \cdot \text{m}^2 / \text{C}^2 = \text{Coulomb's constant}$$

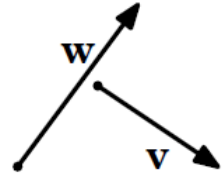
The electrostatic force had the same functional form as Newton's law of gravity

The magnitude of the electrostatic force between two-point charges:

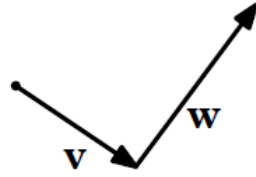
- 1) directly proportional to the product of the magnitudes of charges
- 2) inversely proportional to the square of the distance between them
- 3) The force is along the straight line joining them.

Vector calculus:

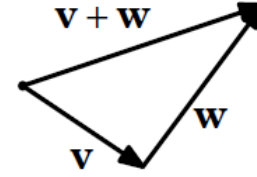
$$\mathbf{v} + \mathbf{w}$$



(a) Vectors \mathbf{v} and \mathbf{w}



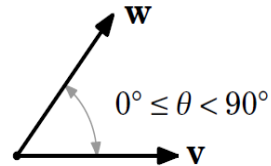
(b) Translate \mathbf{w} to the end of \mathbf{v}



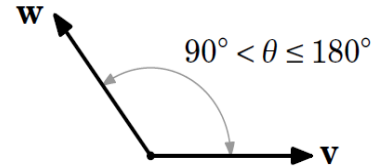
(c) The sum $\mathbf{v} + \mathbf{w}$

Dot product

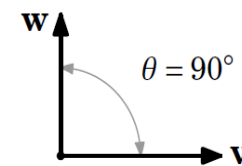
$$\mathbf{v} \cdot \mathbf{w} = \cos \theta \|\mathbf{v}\| \|\mathbf{w}\|$$



(a) $\mathbf{v} \cdot \mathbf{w} > 0$



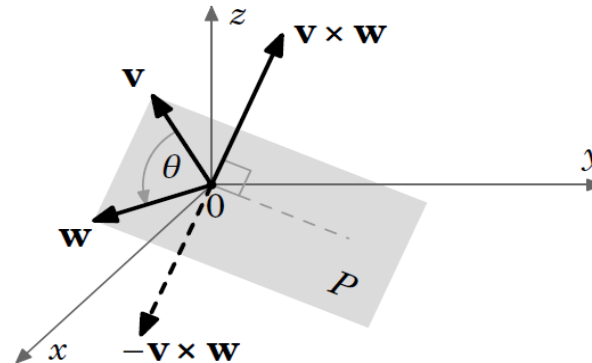
(b) $\mathbf{v} \cdot \mathbf{w} < 0$



(c) $\mathbf{v} \cdot \mathbf{w} = 0$

cross product

$$\|\mathbf{v} \times \mathbf{w}\| = \|\mathbf{v}\| \|\mathbf{w}\| \sin \theta$$



Vector force:

Electric forces follow the law of superposition.

If more than one charge is causing a force on object 1, then the net force acting on object 1 is just the sum of all the individual forces acting on 1.

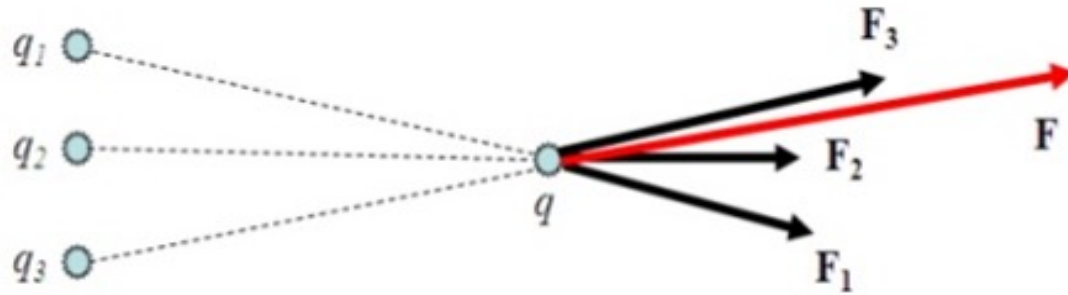
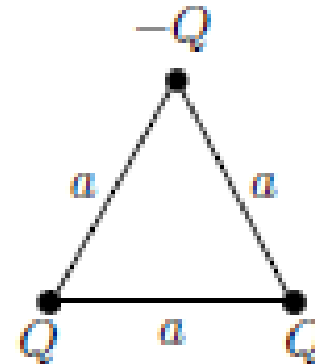


Fig. Superposition Law

$$\text{Net force on } q: \mathbf{F} = \mathbf{F}_1 + \mathbf{F}_2 + \mathbf{F}_3$$

$$\vec{\mathbf{F}}_{tot} = \sum_{i=1}^n \frac{qq_i}{4\pi\epsilon_0 r_i^2} \hat{\mathbf{r}}_i$$



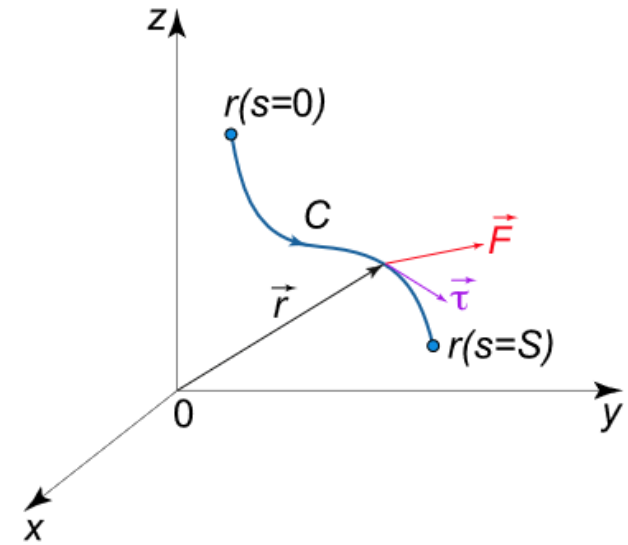
Line integral of vector

Line integral of vector force \mathbf{F} along the curve C .

Suppose that a curve C is defined by the vector function $\mathbf{r}=\mathbf{r}(s)$, $0\leq s\leq S$, where s is the arc length of the curve. Then the derivative of the vector function

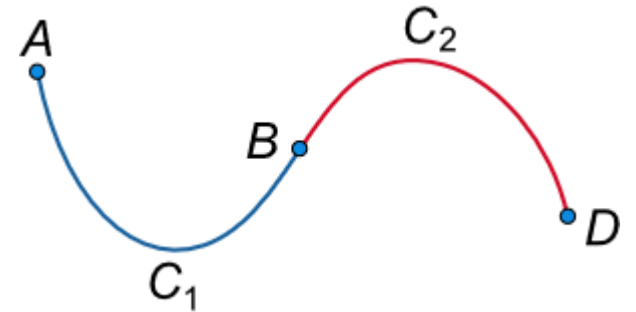
$$\frac{d\mathbf{r}}{ds} = \boldsymbol{\tau} \text{ (Tangent direction at each point of the curve)}$$

Curve direction is important.



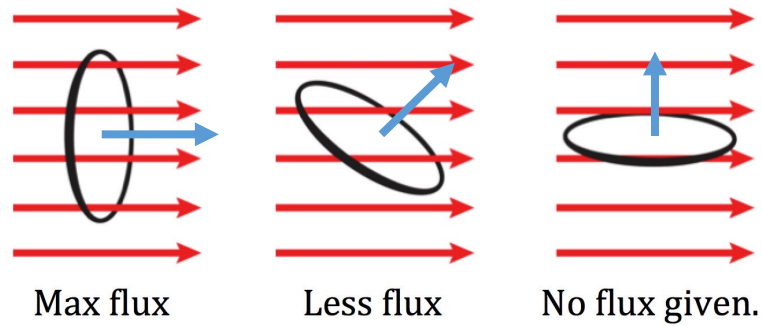
$$\int_C (\mathbf{F} \cdot d\mathbf{r}) = \int_0^S (\mathbf{F}(\mathbf{r}(s)) \cdot \boldsymbol{\tau}) ds,$$

$$\int_C (\mathbf{F} \cdot d\mathbf{r}) = \int_{C_1 \cup C_2} (\mathbf{F} \cdot d\mathbf{r}) = \int_{C_1} (\mathbf{F} \cdot d\mathbf{r}) + \int_{C_2} (\mathbf{F} \cdot d\mathbf{r});$$

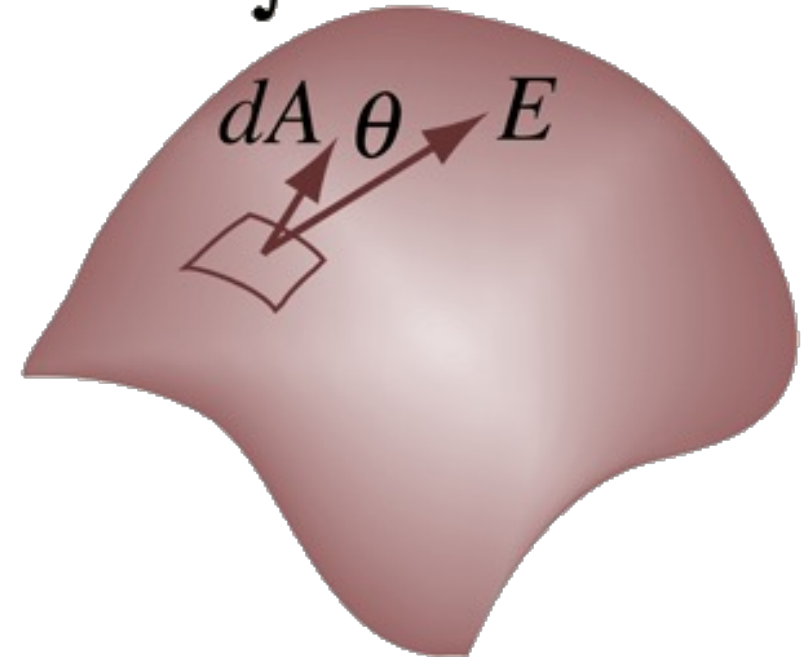


Surface integral of vector

dA direction is perpendicular to the tangent plane to that surface at A



$$\int \vec{E} \cdot d\vec{A} = \int E \cos \theta dA$$



$$\iint_S \vec{F} \cdot d\vec{S} = \iint_S \vec{F} \cdot \vec{n} dS$$

Gradient: Greatest rate of increase

Gradient: 3-dimension derivative of a scalar function

showing the **direction and rate of fastest increase** of the scalar function f at a point space.

How quickly something changes from one point to another

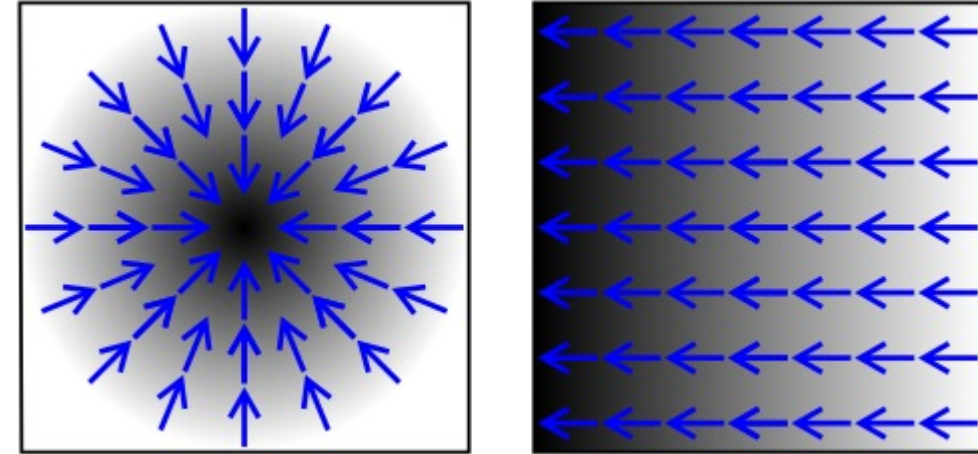
$$\nabla f = \text{grad } f = \left\langle \frac{\partial f}{\partial x}(x, y, z), \frac{\partial f}{\partial y}(x, y, z), \frac{\partial f}{\partial z}(x, y, z) \right\rangle$$

$$\nabla f = \frac{\partial f}{\partial x} \hat{x} + \frac{\partial f}{\partial y} \hat{y} + \frac{\partial f}{\partial z} \hat{z}$$

f : Scalar function

∇f (Gradient): Vector function

Direction: fastest rate of increase

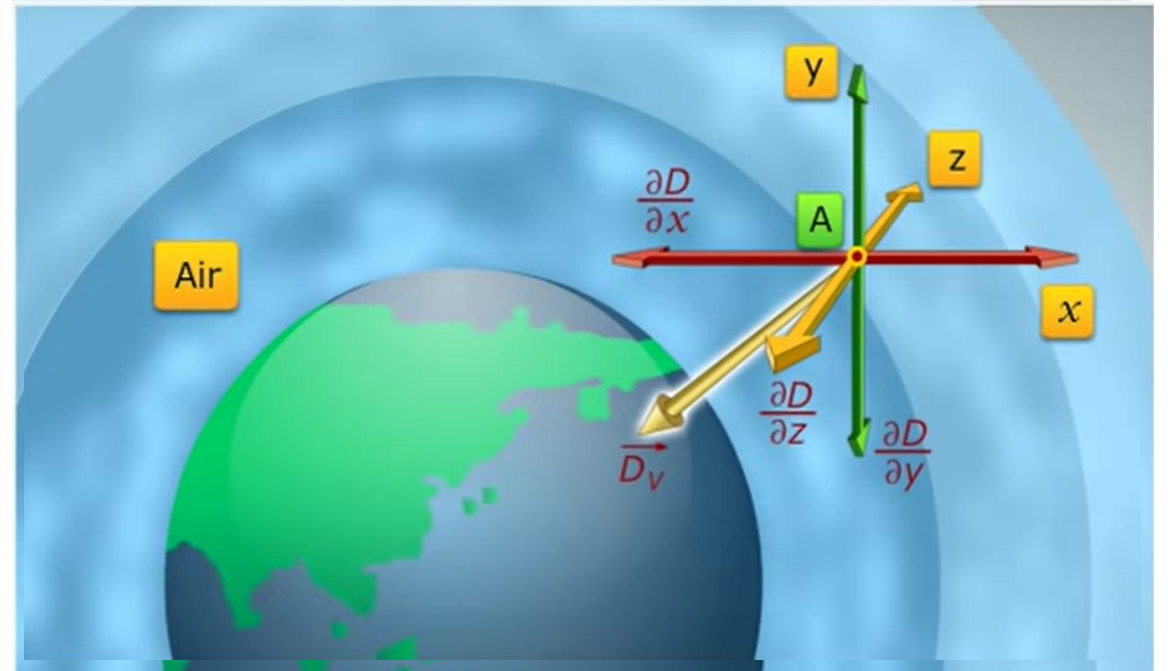


Example: air density in the space

$$D = f(x, y, z)$$

$$\vec{D}_V = \frac{\partial D}{\partial x} \hat{i} + \frac{\partial D}{\partial y} \hat{j} + \frac{\partial D}{\partial z} \hat{k}$$

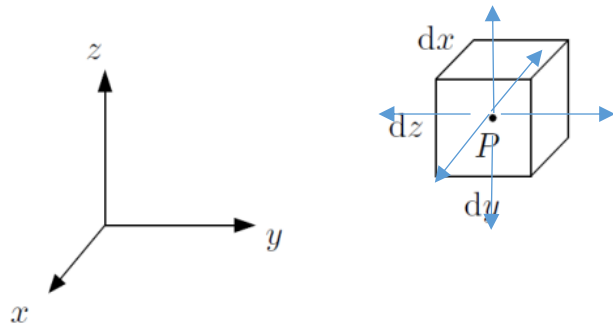
$|\vec{D}_V|$ = **Maximum Rate** at which the **Density Increases**



Divergence: Flux/field out of a point

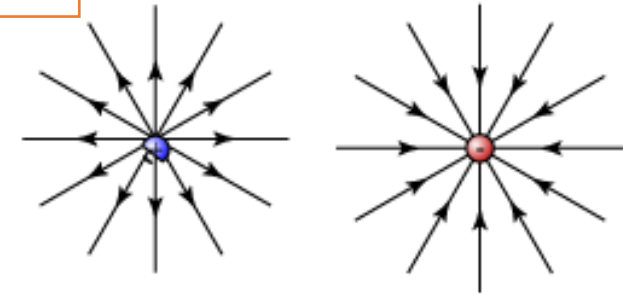
Divergence represents the volume density of the outward flux of a vector field from an infinitesimal volume around a given point

$$\operatorname{div} \mathbf{F} = \nabla \cdot \mathbf{F} = \lim_{V \rightarrow 0} \frac{1}{|V|} \oiint_{S(V)} \mathbf{F} \cdot \hat{\mathbf{n}} dS$$

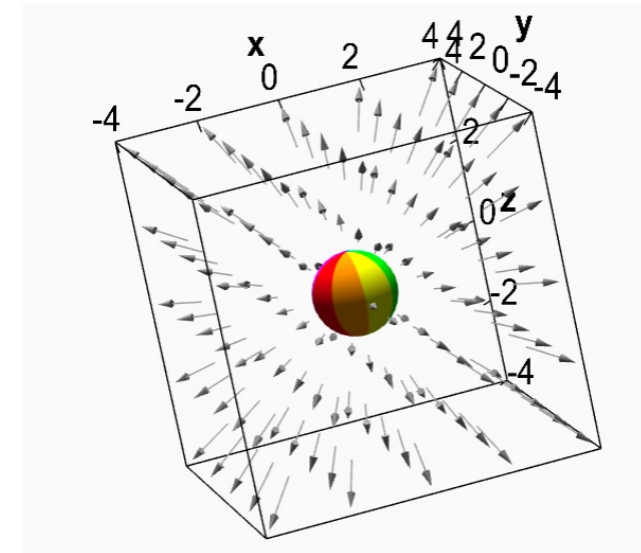


$$\nabla \cdot \mathbf{F} = \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z}$$

\mathbf{F} is a vector
 $\nabla \cdot \mathbf{F}$ is a scalar.

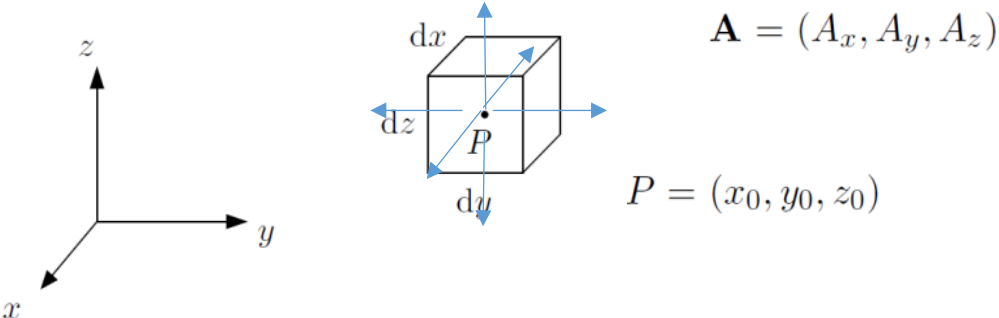


Field pattern of a pointed electrode



Divergence: Mathematical calculation

$$\boxed{\operatorname{div} \mathbf{A} = \lim_{\Delta v \rightarrow 0} \frac{\oint_S \mathbf{A} \cdot d\mathbf{S}}{\Delta v}}$$



$$\begin{aligned} \oint_S \mathbf{A} \cdot d\mathbf{S} &= \oint_S \mathbf{A} \cdot \hat{\mathbf{n}} dS \\ \oint_S \mathbf{A} \cdot d\mathbf{S} &= \int_{\text{foran}} A_x(\text{foran}) dydz - \int_{\text{bak}} A_x(\text{bak}) dydz \quad \leftarrow A_x(\text{foran}) - A_x(\text{bak}) = A_x(x_0 + dx/2, y_0, z_0) - A_x(x_0 - dx/2, y_0, z_0) = \frac{\partial A_x}{\partial x} dx \\ &\quad - \int_{\text{venstre}} A_y(\text{venstre}) dx dz + \int_{\text{høyre}} A_y(\text{høyre}) dx dz \quad \leftarrow A_y(\text{høyre}) - A_y(\text{venstre}) = \frac{\partial A_y}{\partial y} dy \\ &\quad + \int_{\text{topp}} A_z(\text{topp}) dx dy - \int_{\text{bunn}} A_z(\text{bunn}) dx dy. \quad \leftarrow A_z(\text{topp}) - A_z(\text{bunn}) = \frac{\partial A_z}{\partial z} dz. \end{aligned}$$



$$\oint_S \mathbf{A} \cdot d\mathbf{S} = \frac{\partial A_x}{\partial x} dx dy dz + \frac{\partial A_y}{\partial y} dx dy dz + \frac{\partial A_z}{\partial z} dx dy dz$$



$$\boxed{\operatorname{div} \mathbf{A} = \lim_{\Delta v \rightarrow 0} \frac{\oint_S \mathbf{A} \cdot d\mathbf{S}}{\Delta v}}$$

$\Delta v = dx dy dz$

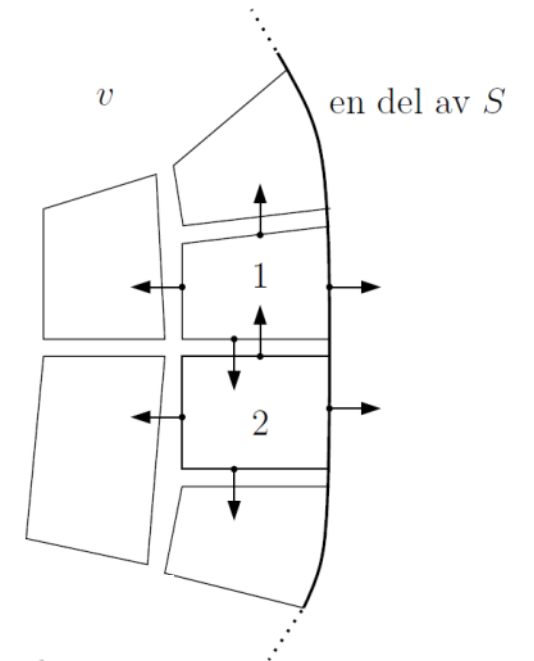


$$\boxed{\operatorname{div} \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} = \nabla \cdot \mathbf{A}}$$

Divergence theorem

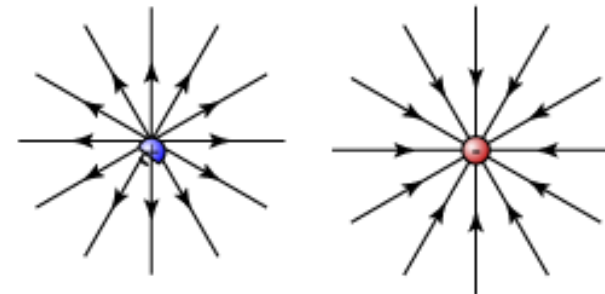
$$\oint_S \mathbf{A} \cdot d\mathbf{S} = \int_v \nabla \cdot \mathbf{A} dv$$

$$\oint_S \mathbf{E} \cdot d\mathbf{S} = \sum_i \oint_{S_i} \mathbf{E} \cdot d\mathbf{S} = \sum_i \left(\frac{1}{\Delta V_i} \oint_{S_i} \mathbf{E} \cdot d\mathbf{S}_i \right) \Delta V_i \rightarrow \int \nabla \cdot \mathbf{E} dV.$$



$$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$$

Gauss' Law



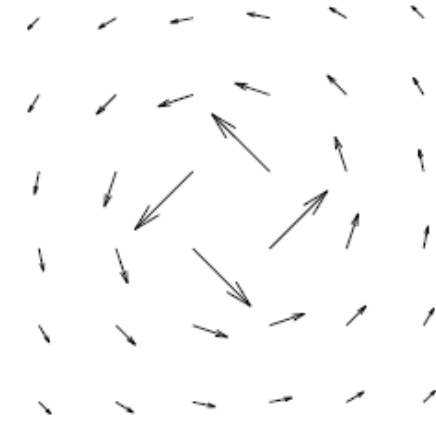
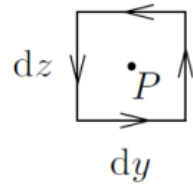
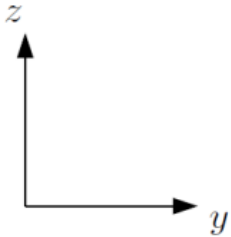
Field pattern of
a pointed electrode

Curl: how much does a field circulate around a point.

The curl of a field presents the infinitesimal circulation density at each point of the field

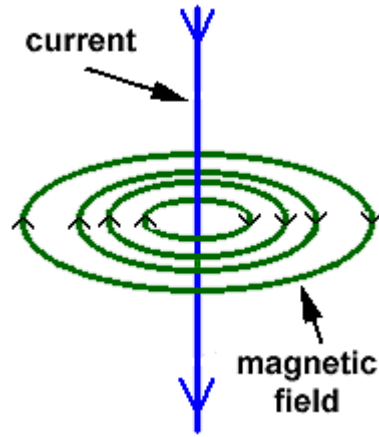
One dimension x

$$(\text{curl } \mathbf{A})_x = \lim_{\Delta S \rightarrow 0} \frac{\oint_C \mathbf{A} \cdot d\mathbf{l}}{\Delta S}$$



sirkulasjon/curl

Three dimensions x, y and z



$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{z}$$

Curl

The curl around x-axis, in yz plane

$$\oint_C \mathbf{A} \cdot d\mathbf{l} = \int_{\text{nede}} A_y \, dy - \int_{\text{oppe}} A_y \, dy - \int_{\text{venstre}} A_z \, dz + \int_{\text{høyre}} A_z \, dz$$

$$A_y(\text{nede}) - A_y(\text{oppe}) = A_y(x_0, y_0, z_0 - dz/2) - A_y(x_0, y_0, z_0 + dz/2) = -\frac{\partial A_y}{\partial z} dz,$$

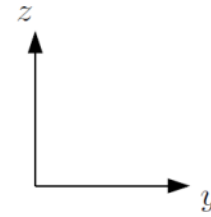
$$A_z(\text{høyre}) - A_z(\text{venstre}) = A_z(x_0, y_0 + dy/2, z_0) - A_z(x_0, y_0 - dy/2, z_0) = \frac{\partial A_z}{\partial y} dy,$$

$$\oint_C \mathbf{A} \cdot d\mathbf{l} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) dy dz,$$

$$(\text{curl } \mathbf{A})_x = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}.$$

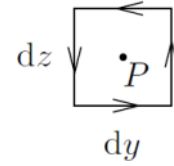
Similar to the curl around y and z-axis

$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{z}$$



$$P = (x_0, y_0, z_0)$$

$$\mathbf{A} = (A_x, A_y, A_z)$$

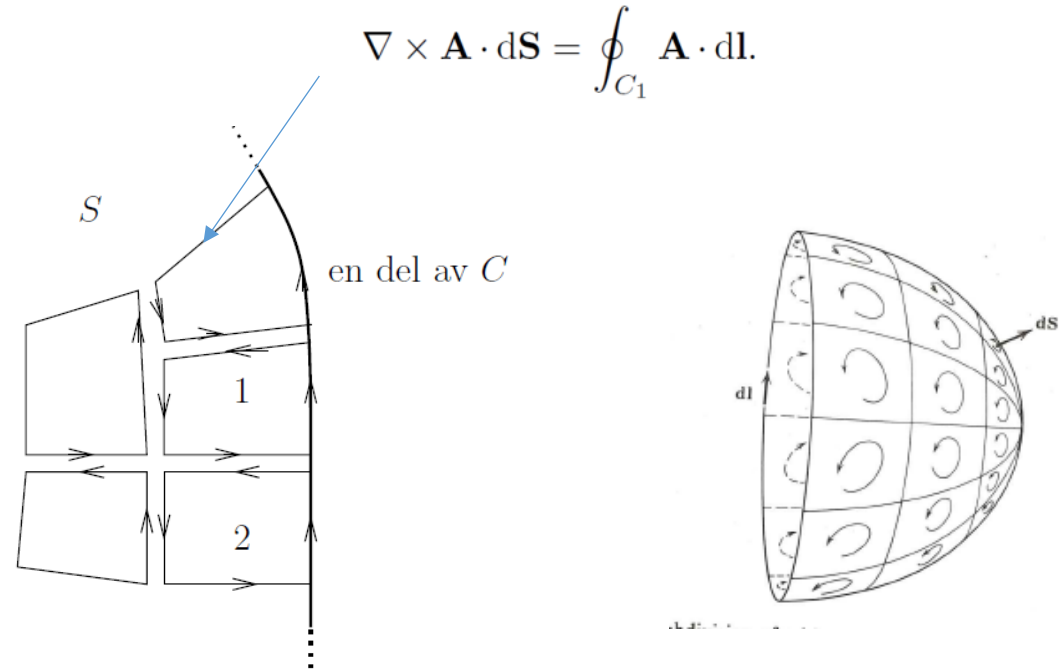


$$\Delta S = dy dz$$

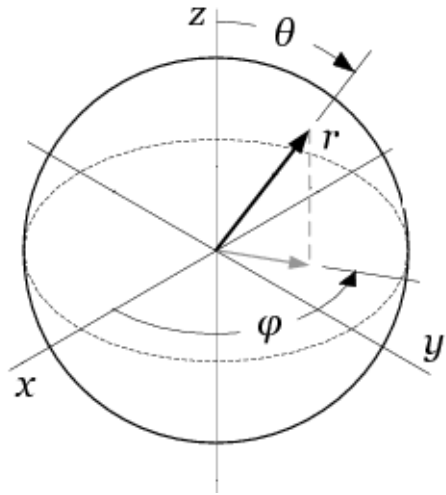
Stokes' Theorem

$$\oint_C \mathbf{A} \cdot d\mathbf{l} = \int_S \nabla \times \mathbf{A} \cdot d\mathbf{S}.$$

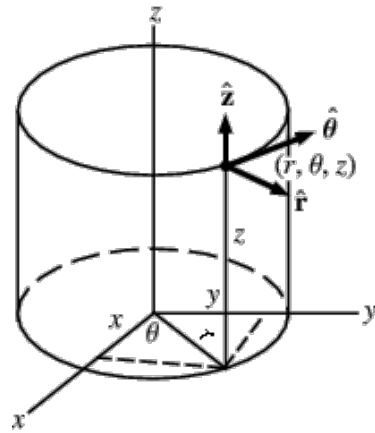
$$\int_S \nabla \times \mathbf{A} \cdot d\mathbf{S} = \sum_i \oint_{C_i} \mathbf{A} \cdot d\mathbf{l}.$$



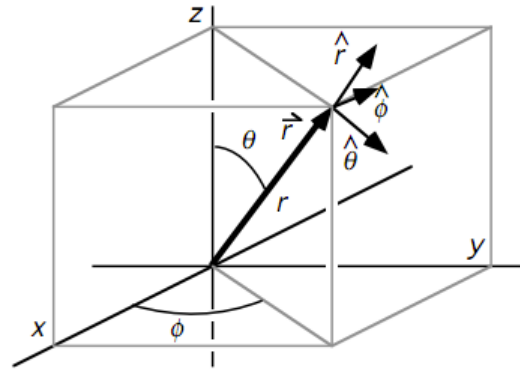
Different coordinates



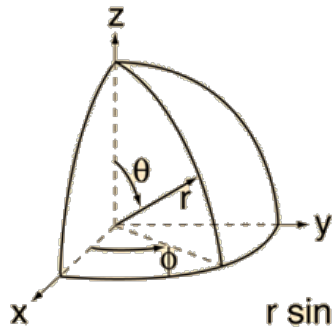
Spherical coordinate



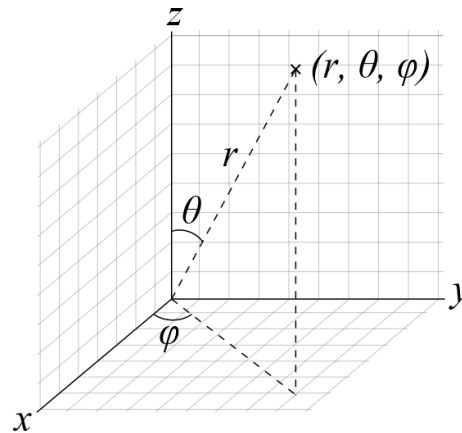
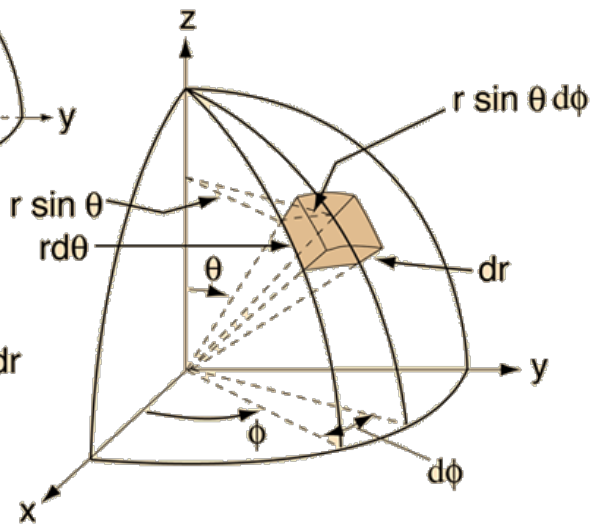
Cylindrical coordinate



Cartesia Coordinate



Volume element
 $dV = r^2 \sin \theta \, d\theta \, d\phi \, dr$



VECTOR DIFFERENTIAL OPERATIONS

Rectangular Coordinates (x, y, z)

$$\begin{aligned} \nabla \Phi &= \hat{x} \frac{\partial \Phi}{\partial x} + \hat{y} \frac{\partial \Phi}{\partial y} + \hat{z} \frac{\partial \Phi}{\partial z} \\ \nabla \cdot \mathbf{D} &= \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z} \\ \nabla \times \mathbf{H} &= \hat{x} \left(\frac{\partial H_z}{\partial y} - \frac{\partial H_y}{\partial z} \right) + \hat{y} \left(\frac{\partial H_x}{\partial z} - \frac{\partial H_z}{\partial x} \right) + \hat{z} \left(\frac{\partial H_y}{\partial x} - \frac{\partial H_x}{\partial y} \right) \\ \nabla^2 \Phi &= \frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2} \\ \nabla^2 \mathbf{A} &= \hat{x} \nabla^2 A_x + \hat{y} \nabla^2 A_y + \hat{z} \nabla^2 A_z \end{aligned}$$

Cylindrical Coordinates (r, ϕ, z)

$$\begin{aligned} \nabla \Phi &= \hat{r} \frac{\partial \Phi}{\partial r} + \hat{\phi} \frac{1}{r} \frac{\partial \Phi}{\partial \phi} + \hat{z} \frac{\partial \Phi}{\partial z} \\ \nabla \cdot \mathbf{D} &= \frac{1}{r} \frac{\partial}{\partial r} (r D_r) + \frac{1}{r} \frac{\partial D_\phi}{\partial \phi} + \frac{\partial D_z}{\partial z} \\ \nabla \times \mathbf{H} &= \hat{r} \left[\frac{1}{r} \frac{\partial H_z}{\partial \phi} - \frac{\partial H_\phi}{\partial z} \right] + \hat{\phi} \left[\frac{\partial H_r}{\partial z} - \frac{\partial H_z}{\partial r} \right] + \hat{z} \left[\frac{1}{r} \frac{\partial (r H_\phi)}{\partial r} - \frac{1}{r} \frac{\partial H_r}{\partial \phi} \right] \\ \nabla^2 \Phi &= \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \phi^2} + \frac{\partial^2 \Phi}{\partial z^2} \\ \nabla^2 \mathbf{A} &= \hat{r} \left(\nabla^2 A_r - \frac{2}{r^2} \frac{\partial A_\phi}{\partial \phi} - \frac{A_r}{r^2} \right) + \hat{\phi} \left(\nabla^2 A_\phi + \frac{2}{r^2} \frac{\partial A_r}{\partial \phi} - \frac{A_\phi}{r^2} \right) + \hat{z} (\nabla^2 A_z) \end{aligned}$$

Spherical Coordinates (r, θ, ϕ)

$$\begin{aligned} \nabla \Phi &= \hat{r} \frac{\partial \Phi}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial \Phi}{\partial \theta} + \frac{\hat{\phi}}{r \sin \theta} \frac{\partial \Phi}{\partial \phi} \\ \nabla \cdot \mathbf{D} &= \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 D_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta D_\theta) + \frac{1}{r \sin \theta} \frac{\partial D_\phi}{\partial \phi} \\ \nabla \times \mathbf{H} &= \frac{\hat{r}}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (H_\phi \sin \theta) - \frac{\partial H_\theta}{\partial \phi} \right] \\ &\quad + \frac{\hat{\theta}}{r} \left[\frac{1}{\sin \theta} \frac{\partial H_r}{\partial \phi} - \frac{\partial}{\partial r} (r H_\phi) \right] + \frac{\hat{\phi}}{r} \left[\frac{\partial}{\partial r} (r H_\theta) - \frac{\partial H_r}{\partial \theta} \right] \\ \nabla^2 \Phi &= \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \Phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2} \\ \nabla^2 \mathbf{A} &= \hat{r} \left[\nabla^2 A_r - \frac{2}{r^2} \left(A_r + \cot \theta A_\theta + \csc \theta \frac{\partial A_\phi}{\partial \phi} + \frac{\partial A_\theta}{\partial \theta} \right) \right] \\ &\quad + \hat{\theta} \left[\nabla^2 A_\theta - \frac{1}{r^2} \left(\csc^2 \theta A_\theta - 2 \frac{\partial A_r}{\partial \theta} + 2 \cot \theta \csc \theta \frac{\partial A_\phi}{\partial \phi} \right) \right] \\ &\quad + \hat{\phi} \left[\nabla^2 A_\phi - \frac{1}{r^2} \left(\csc^2 \theta A_\phi - 2 \csc \theta \frac{\partial A_r}{\partial \phi} - 2 \cot \theta \csc \theta \frac{\partial A_\theta}{\partial \phi} \right) \right] \end{aligned}$$

Examples: Problem 3

Calculate the integral

$$I = \int_V (\nabla \cdot \mathbf{F}) dV \quad (1)$$

where $\mathbf{F} = r\hat{\mathbf{r}} = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$, and the volume V is a sphere with radius R placed in the origin.

- Calculate the integral directly.
- Calculate the integral using the divergence theorem.

$$\operatorname{div} \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} = \nabla \cdot \mathbf{A}.$$

$$\oint_S \mathbf{A} \cdot d\mathbf{S} = \int_V \nabla \cdot \mathbf{A} dv$$

Solutions:

Calculate the integral

$$I = \int_V (\nabla \cdot \mathbf{F}) dV \quad (1)$$

where $\mathbf{F} = r\hat{\mathbf{r}} = x\hat{\mathbf{x}} + y\hat{\mathbf{y}} + z\hat{\mathbf{z}}$, and the volume V is a sphere with radius R placed in the origin.

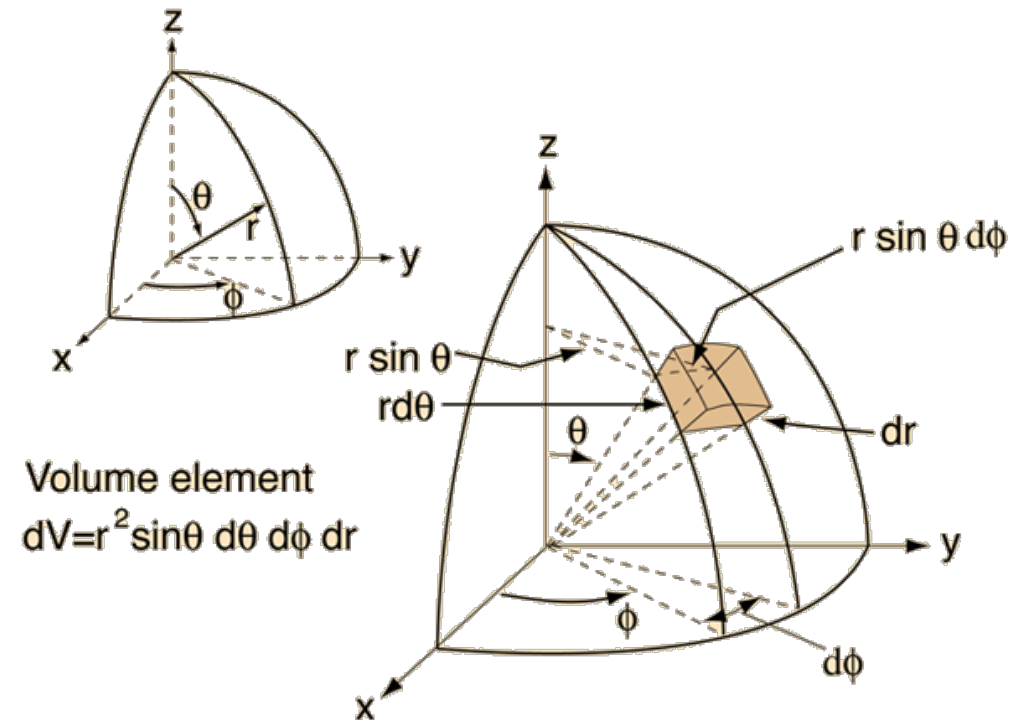
- Calculate the integral directly.
- Calculate the integral using the divergence theorem.

$$\text{div} \mathbf{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z} = \nabla \cdot \mathbf{A}.$$

$$\nabla \cdot \mathbf{F} = 3.$$

$$\begin{aligned} \int_V (\nabla \cdot \mathbf{F}) dV &= \int_0^R \int_0^{2\pi} \int_0^\pi 3r^2 \sin \theta d\varphi d\theta dr \\ &= 3 \int_0^R r^2 dr \int_0^\pi \sin \theta d\theta \int_0^{2\pi} d\varphi \\ &= \underline{\underline{4\pi R^3}}. \end{aligned}$$

$$\underline{\underline{\int_V (\nabla \cdot \mathbf{F}) dV = 3 \int_V dV = 4\pi R^3.}}$$



Solution for b

$$\begin{aligned}\int_v (\nabla \cdot \mathbf{F}) dv &= \oint_S \mathbf{F} \cdot d\mathbf{S} \\ &= \int_0^{2\pi} \int_0^\pi (R\hat{\mathbf{r}}) \cdot (R^2 \sin\theta d\theta d\varphi \hat{\mathbf{r}}) \\ &= R^3 \int_0^{2\pi} d\varphi \int_0^\pi \sin\theta d\theta \\ &= \underline{\underline{4\pi R^3}}.\end{aligned}$$

$$\oint_S \mathbf{A} \cdot d\mathbf{S} = \int_v \nabla \cdot \mathbf{A} dv$$

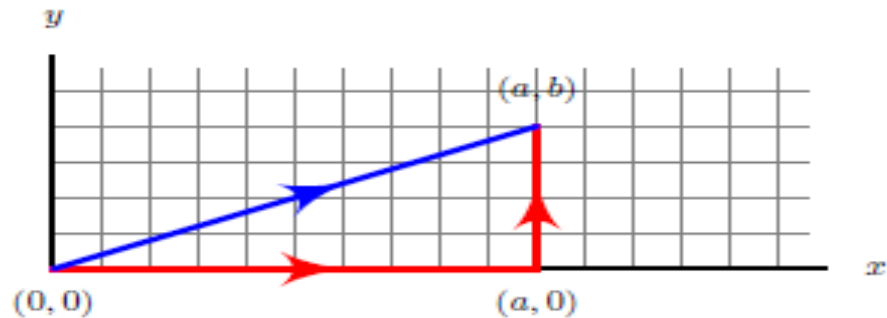
Example:

Calculate the line integral

$$I = \int_C \mathbf{F} \cdot d\mathbf{l}, \quad (2)$$

Where $\mathbf{F} = (xy^2 + 2y)\hat{x} + (x^2y + 2x)\hat{y}$,

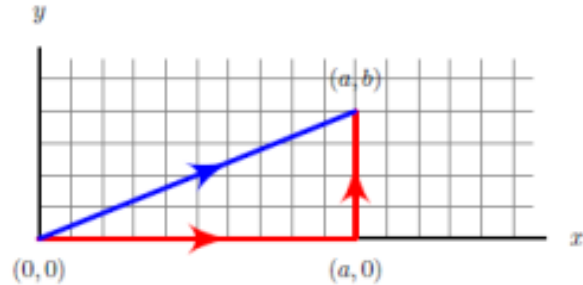
- Along the curve C_1 which consists of two straight lines connecting the points $(0, 0)$, $(a, 0)$ and (a, b) , see figure below.
- Along the curve C_2 which consists of one straight line connecting the points $(0, 0)$ and (a, b) , see figure below.
- Why do these calculations produce the same answer? Explain using Stoke's theorem.



$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{z}$$

$$\oint_C \mathbf{A} \cdot d\mathbf{l} = \int_S \nabla \times \mathbf{A} \cdot d\mathbf{S}.$$

Solution for i) and ii)



$$I = \int_V (\nabla \cdot \mathbf{F}) dV$$

$$\mathbf{F} = (xy^2 + 2y)\hat{\mathbf{x}} + (x^2y + 2x)\hat{\mathbf{y}}$$

- i) Along the curve C_1 which consists of two straight lines connecting the points $(0, 0)$, $(a, 0)$ and (a, b) , see figure below.

$$\begin{aligned} I &= \int_C \mathbf{F} \cdot d\mathbf{l} \\ &= \int_0^b (a^2y + 2a) dy \\ &= \underline{\underline{\frac{1}{2}a^2b^2 + 2ab.}} \end{aligned}$$

- ii) Along the curve C_2 which consists of one straight line connecting the points $(0, 0)$ and (a, b) , see figure below.

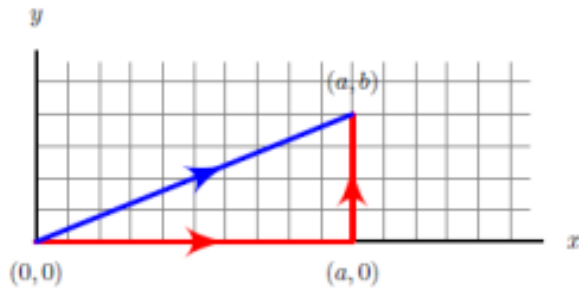
$$d\mathbf{l} = dx\hat{\mathbf{x}} + dy\hat{\mathbf{y}}.$$

$$\begin{aligned} \mathbf{F} \cdot d\mathbf{l} &= F_x dx + F_y dy \\ &= (xy^2 + 2y)dx + (x^2y + 2x)dy. \end{aligned}$$

$$y = \frac{bx}{a}$$

$$\begin{aligned} I &= \int_C \mathbf{F} \cdot d\mathbf{l} \\ &= \int_0^a \left[x \left(\frac{bx}{a} \right)^2 + 2 \left(\frac{bx}{a} \right) \right] dx + \int_0^b \left[\left(\frac{ax}{b} \right)^2 y + 2 \left(\frac{ax}{b} \right) \right] dy \\ &= \underline{\underline{\frac{1}{2}a^2b^2 + 2ab.}} \end{aligned}$$

Conservative vector: solution for iii)



iii) Why do these calculations produce the same answer? Explain using Stoke's theorem.

$$\mathbf{F} = (xy^2 + 2y)\hat{x} + (x^2y + 2x)\hat{y}$$

These integrals have equal values since \mathbf{F} is a conservative field:

$$\begin{aligned}\nabla \times \mathbf{F} &= \left(\frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} \right) \hat{z} \\ &= (2xy - 2xy) \hat{z} \\ &= 0.\end{aligned}$$

$$\nabla \times \mathbf{A} = \left(\frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z} \right) \hat{x} + \left(\frac{\partial A_x}{\partial z} - \frac{\partial A_z}{\partial x} \right) \hat{y} + \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) \hat{z}$$

$$\oint_C \mathbf{A} \cdot d\mathbf{l} = \int_S \nabla \times \mathbf{A} \cdot d\mathbf{S}.$$